

# Hot air motor as cryocooler

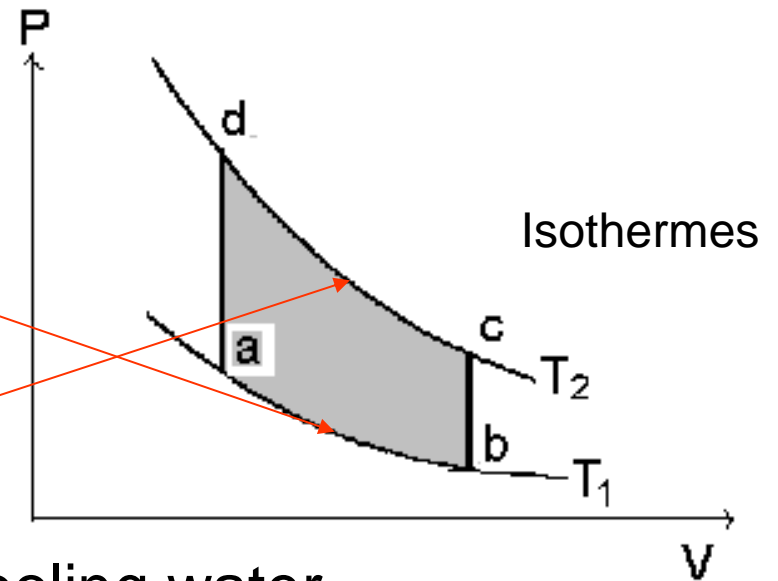
Upper room :Heat will be extracted

← → Cooling water is needed

Process gas above  
a-b:Expansion: Process gas  
detract heat from environment

Heat will be extracted out of copper wool

b-c: Displacing piston up  $T_2$



c-d: Compression → heat goes into cooling water

d-a: Displacing piston down to  $T_1$

Copper wool takes up heat: Gas cools

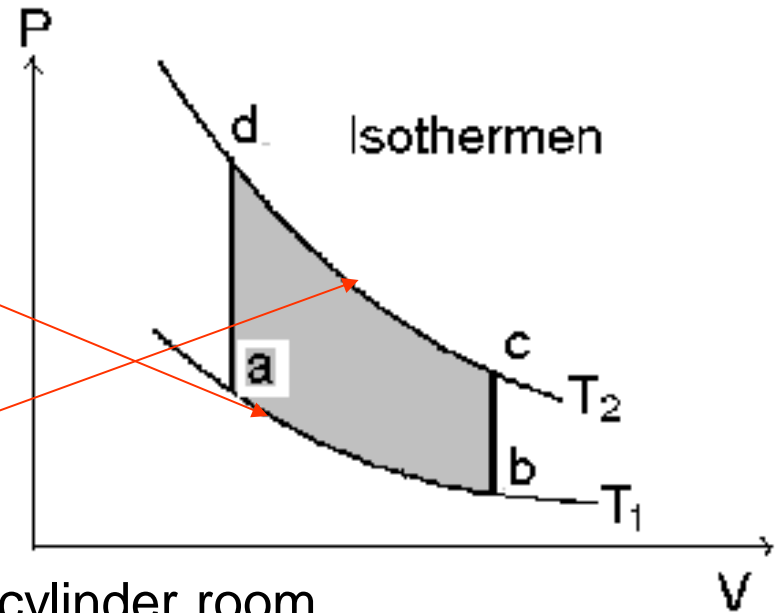
# Hot air motor as heat pump

Lower room: Heat will be detracted and added to upper room

Process gas below:  
**a-b: Expansion:** Heat from cooling water

b-c: Displacing piston down  
 $T_2$  copper wool loses heat

**c-d: Compression** → Heat towards upper cylinder room



# 5.8. Thermodynamic Work

Gas  $(P,V,T), (P,V,T)$  are state variables

1,2,3,4 are equilibrium states

Considered are **quasistatic processes**:

Change of  $P,T,V$  'very slow, very small'

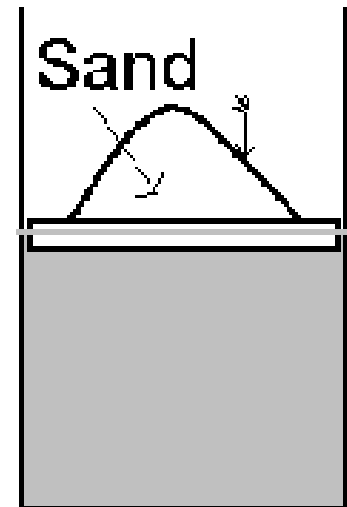
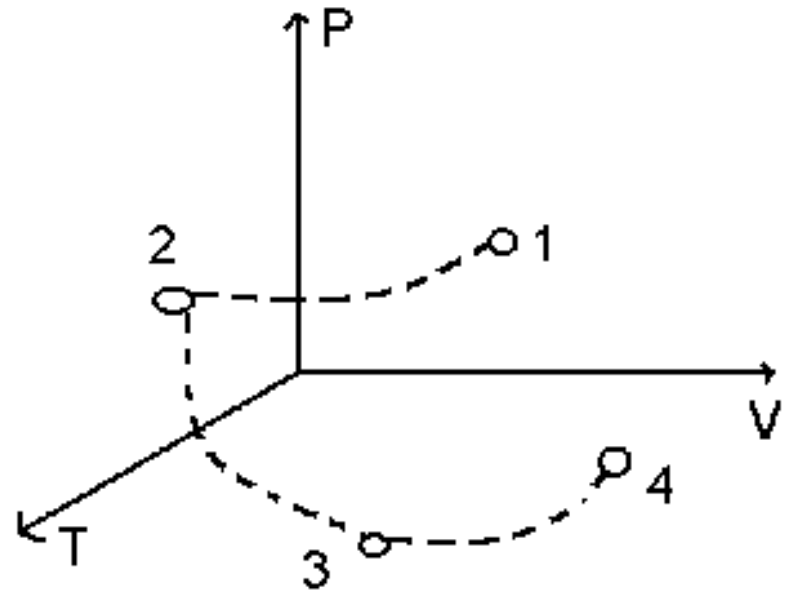
Prozess is reversibel!

e.g: Sand: + grain of sand

**Piston down**

- Sandkörnchen

**Piston back**



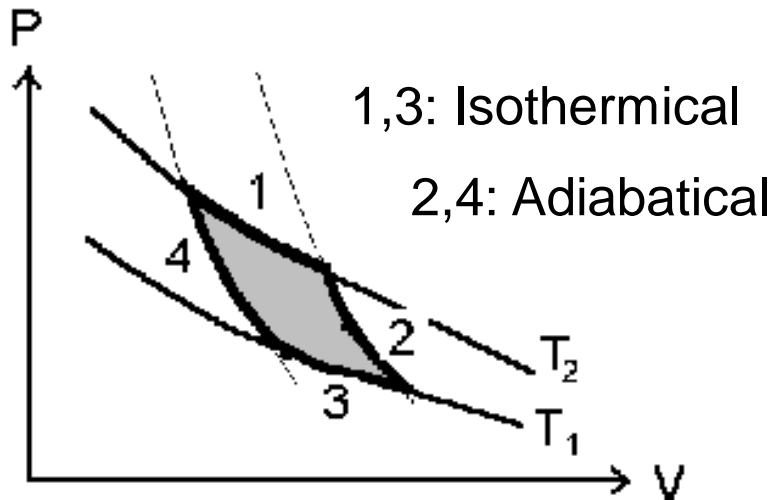
**Irreversible processes:** e.g: Chemical reactions, perfume → air

Work performed:  $W = W(P, V, T)$

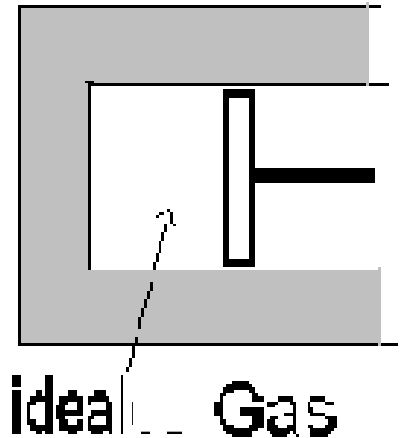
## 5.9. Carnot -cycle

(The most efficient heat cycle)

Exchange of heat possible or interrupted!



Prozesses reversible



$P \cdot V = R \cdot T$	$T \cdot V^{\gamma-1} = \text{constant}$
$U = c_V \cdot T$	$\Delta Q = \Delta U + \Delta W$

$$1: W_1 \text{ positive, } \Delta U = 0 \Rightarrow \Delta Q_{T_2} = W_1$$

$$2: \Delta Q = 0, W_2 \text{ positive,} \\ \Rightarrow \Delta U = W_2 = c_V \cdot (T_2 - T_1)$$

$$3: \Delta U = 0 \Rightarrow W_3 = \Delta Q_{T_1} < 0$$

$$4: \Delta Q = 0 \Rightarrow W_4 = \Delta U = c_V (T_1 - T_2)$$

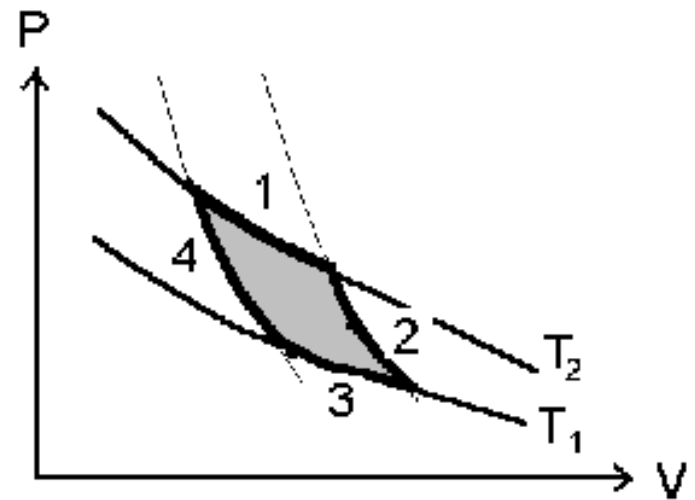
$$\text{Overall: } W = \Delta Q_{T_2} (> 0) + \Delta Q_{T_1} (< 0)$$

### Efficiency of CARNOT- machine

$$W_1 = \int^1 P \cdot dV \quad \text{with} \quad P = \frac{N \cdot k \cdot T}{V} \\ = N \cdot k \cdot T_2 \cdot \int^1 \frac{dV}{V} = N \cdot k \cdot T_2 \cdot \ln \frac{V_2}{V_1};$$

$V_2$  ( Volume after cycle from 1 etc.)

$$\text{analogical: } W_3 = N \cdot k \cdot T_1 \cdot \ln \frac{V_4}{V_3};$$



For adiabatical processes applies:

$$T_2 \cdot V_2^{\gamma-1} = T_1 \cdot V_3^{\gamma-1}$$

$$T_1 \cdot V_4^{\gamma-1} = T_2 \cdot V_1^{\gamma-1}$$

$$\Rightarrow \frac{V_3}{V_4} = \frac{V_2}{V_1} \Rightarrow$$

$$\frac{W_1}{W_3} = \frac{\Delta Q_{T_2}}{\Delta Q_{T_1}} = -\frac{T_2}{T_1}$$

$$\frac{\Delta Q_{T_2}}{T_2} = -\frac{\Delta Q_{T_1}}{T_1} \Rightarrow$$

$$\frac{\Delta Q_{T_2}}{T_2} = -\frac{\Delta Q_{T_1}}{T_1} \Rightarrow$$

$$W = \Delta Q_{T_2} \left(1 - \frac{T_1}{T_2}\right)$$

$$\left(1 - \frac{T_1}{T_2}\right) = \eta \Rightarrow \eta = \frac{\Delta T}{T}$$

Result: Only a part of  $Q_{T_2}$   
can be inverted to work!

Kelvin(1824-1907)

## 5.10. Second fundamental law of thermodynamics

R. Clausius (1822-1888) :

**Not all energy conserving processes are possible**

Heat, without work performed, flows always from higher to lower temperature!

**Clausius** introduced the term **entropy S**.

with  $dQ = T \cdot dS$  or  $dS = \frac{dQ}{T}$

$$\Delta S = S_{end} - S_{beginning} = \int_{beginning}^{end} \frac{dQ}{T} \quad \Delta S_{CARNOT} = 0 : \text{Reversible!}$$

because (s.above):  $\frac{\Delta Q_{T_2}}{T_2} = -\frac{\Delta Q_{T_1}}{T_1}$  Irreversible processes lead to  $\Delta S > 0$

**Clausius: 1. Energy of the world is constant**

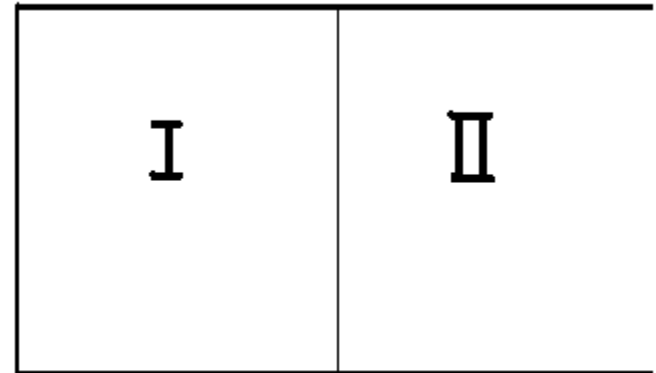
**2. The entropy of the world pursuit to a maximum !"**

# 5.11. Entropy and probability (Boltzmann)

$$N = N_1 + N_2$$

N particles in I and II

Example:



Probability

$N_1$  in I and  $N_2$  in II

e.g.:  $N=8$

One possibility 8 in I  
8 possibilities 1 in I, 7 in II

Uniform distribution:

i.e: 4 in I 4 in II

$$\frac{8!}{4! \cdot 4!} = 70$$

The more P, the smaller  $N_1 - N_2$

P at most:  $N_1 = N_2 = \frac{N}{2}$

At  $N = 10^{23}$

Possibilities for realisation P  
in general:

$$P = \frac{N!}{N_1! \cdot N_2!}$$

All other states are  
improbable!



Irreversible growth of S leads to growing molecular disorder.

Connection between P and S

$$S = k^* \ln P$$

$N_1 = N_2$  in the example above, forms **the endpoint** of an irreversible development!

**Entropy grows into direction of future, not past,  
arrow of time!**

## 5.12. Third fundamental law of thermodynamics

*The thermodynamical state of equilibrium at absolute point sero is a state of maximal order, which only one realisation process  $P=1$  !*