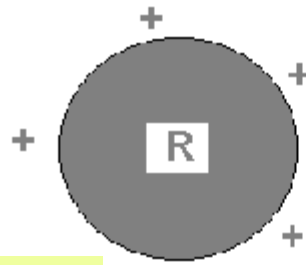


6.6. Capacity



C depends only from shape of the conductor:

Ball with radius R

$$C = 4\pi\epsilon_0 \cdot R$$

Capacitor:

At a given voltage:
Seperation of charge depends on C !

Ball with charge Q against
 ∞

far away wall:

$$U = \frac{1}{4\pi\epsilon_0} \frac{Q}{R}; U \sim Q$$

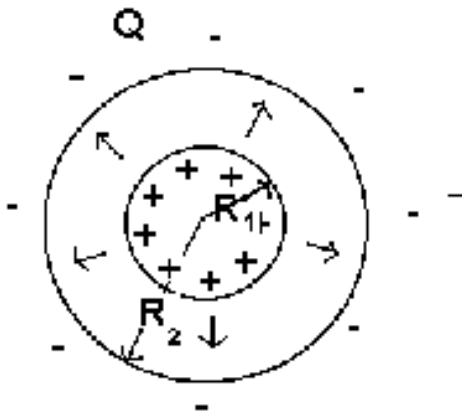
$$Q = C \cdot U,$$

with C as capacity

$$[C] = \text{Farad} = \frac{C}{V}$$

Spherical condensor:

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{Q}{r^2} \vec{e}_r; \quad \text{Potential: } \varphi(r) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r}$$



Voltage U between both ballcups:

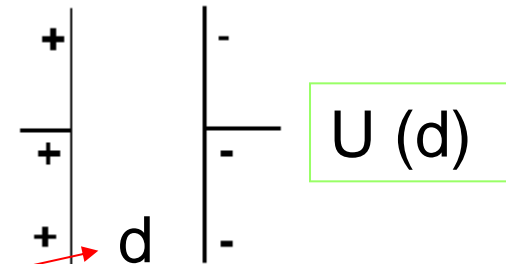
$$U = \varphi_1(R_1) - \varphi_1(R_2) = \frac{1}{4\pi\epsilon_0} \left\{ \frac{Q}{R_1} - \frac{Q}{R_2} \right\}$$

$$= \frac{Q}{4\pi\epsilon_0} \frac{R_2 - R_1}{R_1 \cdot R_2} = \frac{Q}{C} \quad \text{mit} \quad \frac{1}{C} = \frac{1}{4\pi\epsilon_0} \frac{R_2 - R_1}{R_1 \cdot R_2}$$

Special case: $R_1 \approx R_2 = R$ and thus $\Delta R \ll R$

$$C = \epsilon_0 \frac{4\pi R^2}{\Delta R} = \epsilon_0 \frac{A}{\Delta R}$$

Plate condensor:



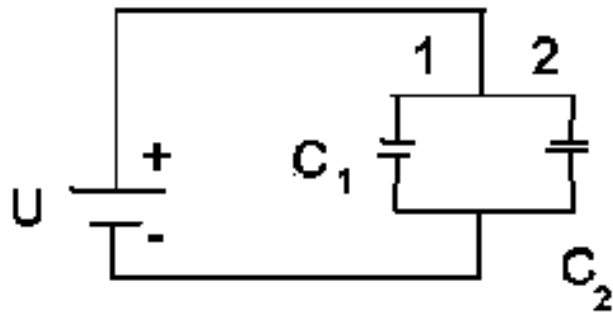
Segment of a spherical condensor with $R \rightarrow \infty$

$$\Rightarrow \boxed{C = \epsilon_0 \frac{A}{d}} \quad d \quad \text{distance of the plates}$$

Potential difference: $U = \frac{Q}{C}$

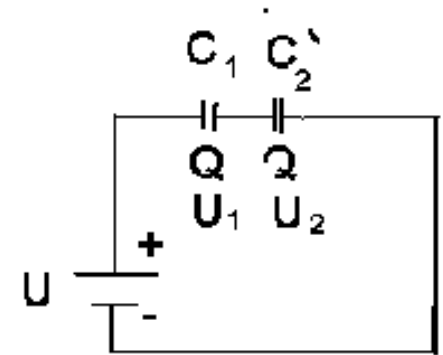
Elektric field: $|\vec{E}| = \frac{d\varphi}{dx} = \frac{\varphi_+ - \varphi_-}{d} = \frac{U}{d} \Rightarrow U = E \cdot d$

Set-up of Condensers



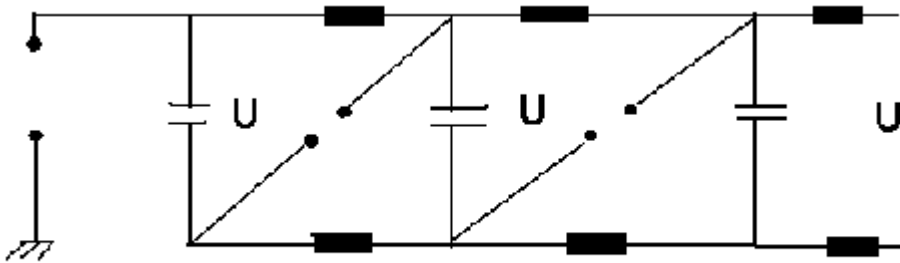
$$Q = Q_1 + Q_2 = C_1 \cdot U + C_2 \cdot U = C \cdot U$$

mit $C_1 + C_2 = C$



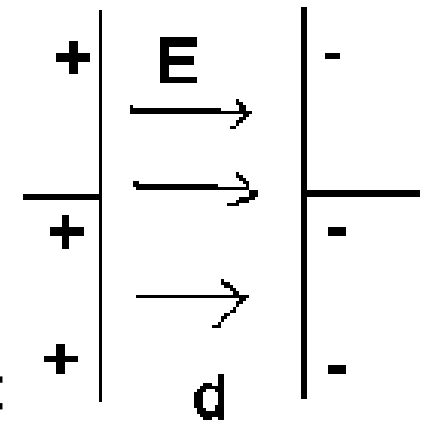
Example: Parallel-serial
Marx-Generator

$$U_1 + U_2 = U; \frac{Q}{C_1} + \frac{Q}{C_2} = \frac{Q}{C}$$



6.7. Energy within a electric field

$$U = \frac{q}{C}; E = \frac{U}{d}$$



Continuing charging $dq \Rightarrow$ work necessary:

$$dW = F \cdot d = dq \cdot E \cdot d = U \cdot dq$$

$$W = \int_0^Q U \cdot dq = \int_0^Q \frac{q}{C} dq = \frac{1}{C} \int_0^Q q \cdot dq = \frac{1}{2} \frac{Q^2}{C} \quad \text{or}$$

$$W = \frac{1}{2} C \cdot U_0^2$$

with $U_0 = \frac{Q}{C}$

Example: $U_0 = 1000V$

$$C = \mu F(\text{arad}) \Rightarrow W = 0.5J = 0.5Ws$$

1Ws:Wattsecond

Where is energy situated?

$$W = \frac{1}{2} C \cdot U_0^2 = \frac{1}{2} \epsilon_0 \frac{A}{d} E_0^2 \cdot d^2 = \frac{1}{2} \epsilon_0 \cdot E_0^2 \cdot V$$

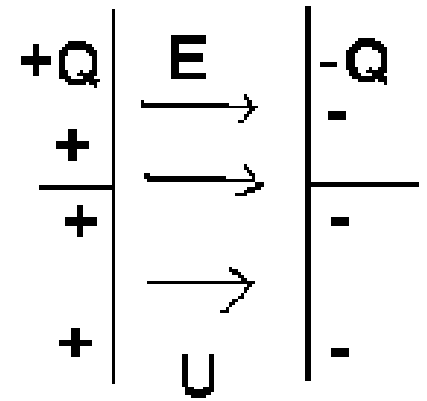
V: Volume of the field region!

Energy is situated in the electric field

Definition of **displacement density**:

$$Q = C \cdot U = \epsilon_0 \frac{A}{d} E \cdot d = \epsilon_0 \cdot A \cdot E$$

Surface density: $\sigma = \frac{Q}{A} = \epsilon_0 \cdot E = D_0$

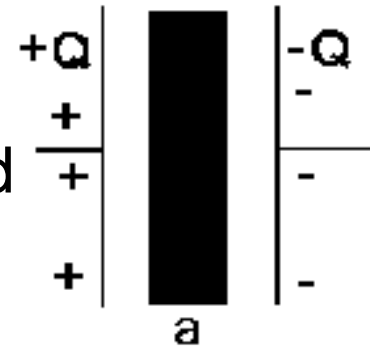


$\vec{D}_0 = \epsilon_0 \cdot \vec{E}$ Displacement density of the vacuum

6.8. Matter in a electric field

Matter: Conductor, distance of the plates: d

$$Q = C_0 \cdot U_0; U_0 = E \cdot d$$



Bringing in of an conductor: Voltage decreases

with conductor: $U = E(d - a) \Rightarrow U = \frac{d-a}{d} \cdot U_0$

because $Q = \text{constant}$ $C = \frac{d}{d-a} \cdot C_0$

i.e.: Capacity increases!

Insulator in a electric field!

Without dielectric matter: $Q = C_0 \cdot U_0$

with: $U = \frac{Q}{C} = \frac{C_0 \cdot U_0}{C} = \frac{U_0}{\epsilon}$

Dielectric constant

$$\epsilon = \frac{C}{C_0}$$

$C = \epsilon \cdot \epsilon_0 \cdot \frac{A}{d}$ for plate condensor

Glas	$\epsilon = 5-10$
Water	81.1 (18°)
Air	1.000576

How to explain it?

Without dielectricum: $\Phi_0 = E_0 \cdot A = \frac{Q_0}{\epsilon_0} : *$

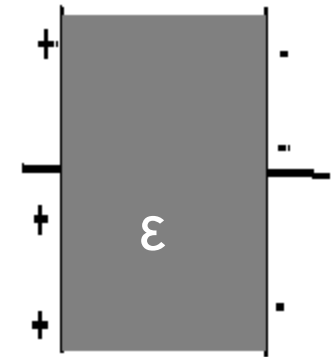
With dielectricum: $U = \frac{U_0}{\epsilon} \Rightarrow E = \frac{E_0}{\epsilon}$

$$\Phi = E \cdot A = \frac{Q}{\epsilon_0} \quad \Phi = \frac{E_0}{\epsilon} \cdot A = \frac{Q_0}{\epsilon_0 \cdot \epsilon} \Rightarrow Q = \frac{Q_0}{\epsilon}$$

On the plates $Q_0!$

Difference Q_P :

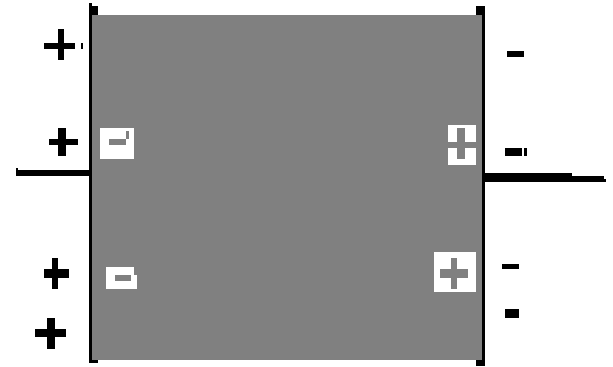
Field lines stop at the surface of the dielectricum



$$Q_P = Q_0 - Q = Q_0 \left(1 - \frac{1}{\epsilon}\right) = \epsilon_0 \cdot \frac{\epsilon - 1}{\epsilon} A \cdot E_0;$$

Surface charge density on the dielectricum

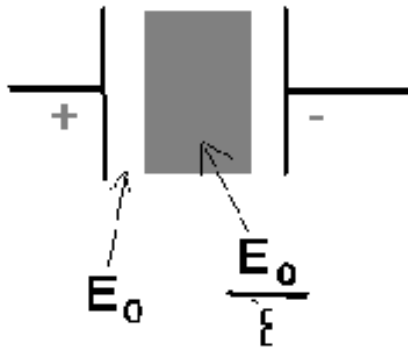
$$\sigma_P = \frac{Q_P}{A} = \epsilon_0 \cdot \frac{\epsilon - 1}{\epsilon} \cdot E_0 = \epsilon_0 \cdot (\epsilon - 1) \cdot E$$



Elektric dipolmoment of the isolator:

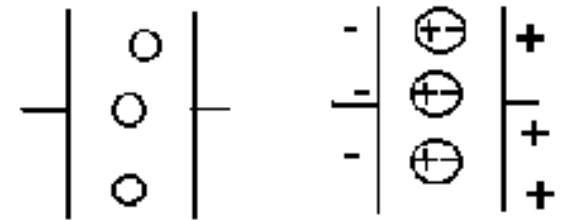
$$\vec{P} = Q_P \cdot \vec{d} = \sigma_P \cdot A \cdot \vec{d} = \epsilon_0 \cdot (\epsilon - 1) \cdot V \cdot \vec{E}$$

Polarisation: $\vec{p} = \frac{\vec{P}}{V}$ $\vec{p} = \epsilon_0 \cdot (\epsilon - 1) \cdot \vec{E} = \epsilon \cdot \epsilon_0 \cdot \vec{E} - \epsilon_0 \cdot \vec{E}$



Polarisation:

a) Molecules without permanent dipolmoment

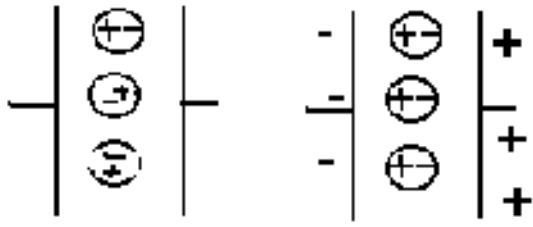


Without field

With field

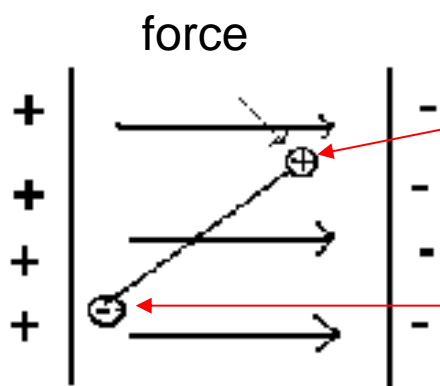
Electric dipolmoments get **induced!**
 ',Displacement polarisation'

b) Molecules with a permanent dipole moment



without field partial polarised

"Orientation polarisation"



force

force: $\vec{F} = Q \cdot \vec{E};$

l : distance of charges

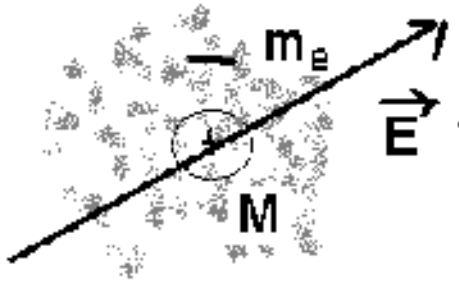
$\vec{F}' = -Q \cdot \vec{E};$

\Rightarrow torque : $\vec{T} = \vec{l} \times \vec{F} = Q \cdot \vec{l} \times \vec{E}$

with : $Q \cdot \vec{l} = \vec{P}$: electric dipole moment

$$\vec{T} = \vec{P} \times \vec{E}$$

Polarisability of atoms in electric alternating fields :



\vec{E} : outer field, masses M, m_e
with $M \gg m_e$

the electron gets moved by force
out of balance

$$m_e \frac{d^2x}{dt^2} + m_e \cdot \omega_0^2 \cdot x = q \cdot E_x^0 \cdot \cos \omega t \quad \text{solution with } x = x_0 \cdot \cos \omega t$$

$$\omega^2 \cdot m_e \cdot x_0 \cdot \cos \omega t + m_e \cdot \omega_0^2 \cdot x_0 \cdot \cos \omega t = q \cdot E_x^0 \cdot \cos \omega t$$

$$x_0 = \frac{q \cdot E_x^0}{m_2(\omega_0^2 - \omega^2)}$$

The deflection corresponds to an oscillating dipole moment

$$p_x = q \cdot x$$

$$m_e = m_2$$

$$\Rightarrow p_x = \frac{q^2 \cdot E_x}{m_2(\omega_0^2 - \omega^2)} = \epsilon_0 \cdot \alpha(\omega) \cdot E_x$$

